

The Newtonian limit: slow motion in weak fields

In this limit all fields are weak and all velocities are small. Only the $_{00}$ component of the energy-momentum tensor is non-vanishing, $T_{00} = \mu$ where μ is the mass density of the matter. Therefore we shall only consider the $_{00}$ component of the Einstein's equation, $R_{00} = \kappa(T_{00} - \frac{1}{2}g_{00}T)$.

In the Newtonian limit $g_{00} = 1 + 2\phi$, $\Gamma_{00}^\alpha = -\phi_{,\alpha}$, $R_{00} = -\phi_{,\alpha}^\alpha = \Delta\phi$, where $\alpha = 1, 2, 3$. The Einstein equation thus turns into the Poisson's equation $\Delta\phi = \frac{1}{2}\kappa\mu$ which is equivalent to the Newtonian theory if we put $\kappa = \frac{8\pi G}{c^4}$, where $G = 6.67 \cdot 10^{-11} m^3 kg^{-1} s^{-2}$ is the Newton's gravitational constant.

Schwarzschild metric: a static spherically symmetric solution of vacuum Einstein equations

A general spherically symmetric static metric can be written as

$$ds^2 = A dt^2 - B dr^2 - r^2(d\theta^2 + \sin^2\theta d\phi^2), \quad (1)$$

where A and B are functions of only the "radius" r . The Christoffel symbols $\Gamma_{bc}^a = \frac{1}{2}g^{ad}(g_{db,c} - g_{bc,d} + g_{cd,b})$ are given as

$$\Gamma_{rr}^r = \frac{1}{2} \frac{B'}{B}, \quad \Gamma_{tr}^t = \frac{1}{2} \frac{A'}{A}, \quad \Gamma_{tt}^r = \frac{1}{2} \frac{A'}{B}, \quad \Gamma_{\theta r}^\theta = \frac{1}{r}, \quad \Gamma_{\theta\theta}^r = -\frac{r}{B}, \quad (2)$$

$$\Gamma_{\phi r}^\phi = \frac{1}{r}, \quad \Gamma_{\phi\phi}^r = -\frac{r \sin^2\theta}{B}, \quad \Gamma_{\phi\theta}^\phi = \cot\theta, \quad \Gamma_{\phi\phi}^\theta = -\sin\theta \cos\theta. \quad (3)$$

The Ricci tensor $R_{ab} = R_{acb}^a$, where $R_{bcd}^a = \Gamma_{bd,c}^a - \Gamma_{bc,d}^a + \Gamma_{bd}^e \Gamma_{ec}^a - \Gamma_{bc}^e \Gamma_{ed}^a$ is then equal

$$R_{tt} = \frac{A''}{2B} + \frac{A'}{B} \left(\frac{1}{r} - \frac{B'}{4B} - \frac{A'}{4A} \right), \quad R_{\theta\theta} = 1 - \left(\frac{r}{B} \right)' - \frac{1}{2} \left(\frac{A'}{A} + \frac{B'}{B} \right) \frac{r}{B}, \quad R_{rr} = -\frac{A''}{2A} + \frac{A'B'}{4AB} + \frac{A'^2}{4A^2} + \frac{B'}{rB}, \quad (4)$$

The vacuum Einstein equations are $R_{ab} = 0$. Making a linear combination $BR_{tt} + AR_{rr} = 0$ we find $A'B + AB' = 0 \Rightarrow \frac{A'}{A} + \frac{B'}{B} = 0 \Rightarrow AB = 1$. From $R_{\theta\theta} = 0$ we then find $\frac{r}{B} = r - R$, where R is an integration constant, and $B = \frac{1}{1-\frac{R}{r}}$, $A = 1 - \frac{R}{r}$. Finally the famous Schwarzschild metric is

$$ds^2 = \left(1 - \frac{R}{r}\right) dt^2 - \left(\frac{1}{1 - \frac{R}{r}}\right) dr^2 - r^2(d\theta^2 + \sin^2\theta d\phi^2). \quad (5)$$

The integration constant R is determined from the Newtonian limit, $R = 2GM$, where M is the mass of the central body. It is called gravitational radius, or Schwarzschild radius.

Exercises

1. Calculate the energy-momentum tensor T_{ab} for a particle with mass m . The action is $S = -m \int ds$.
2. For the metric $ds^2 = \left(1 - \frac{R}{r}\right) dt^2 - \frac{1}{1-\frac{R}{r}} dr^2 - r^2(d\theta^2 + \sin^2\theta d\phi^2)$ write down the geodesic equations for a massive body and for a ray of light.