

### Space-time of an isotropic universe. Cosmological redshift. Hubble constant.

The Friedman's equation for the isotropic universe,  $\frac{3}{a^2}(a^2 + a'^2) = \kappa\epsilon$ , and the energy conservation equation,  $d\epsilon = -(\epsilon + p)3\frac{da}{a}$ , can be integrated for the matter dominated universe, where the pressure is zero,  $p = 0$ , and the energy density  $\epsilon$  is equal to the mass density  $\mu$ ,  $\epsilon = \mu$ . The energy conservation gives  $\mu a^3 = \text{Const}$  and the subsequent integration of the Friedmann's equation gives for a closed universe ( $a^2 > 0$ ) a "Big Bang"  $\rightarrow$  "Big Crunch" scenario:

$$a = a_0(1 - \cos(\eta)) , \quad t = a_0(\eta - \sin(\eta)) \quad (1)$$

For the open isotropic universe the Friedmann's equation provides a "Big Bang"  $\rightarrow$  "Expansion Forever" scenario:

$$a = a_0(\cosh(\eta) - 1) , \quad t = a_0(\sinh(\eta) - \eta) \quad (2)$$

For a flat isotropic universe  $ds^2 = dt^2 - b^2(t)(dx^2 + dy^2 + dz^2)$  the scenario is also "Big Bang"  $\rightarrow$  "Expansion Forever" (see the Exercise):  $\mu b^3 = \text{const}, b = \text{const } t^{2/3}$ .

At early stages with high densities the universe was (probably) rather radiation dominated,  $p = \frac{\epsilon}{3}$ . This, however, doesn't save us from the singular point at  $\eta = 0$ . Indeed, we have (for  $\eta \ll 1$ ) :  $\epsilon a^4 = \text{const}$  ,  $a = \text{const} \cdot t^{1/2}$

### Cosmological redshift. Hubble constant.

In an isotropic universe the radial ( $d\theta = d\phi = 0$ ) propagation of light ( $ds^2 = 0$ ) is described by  $\chi = \pm\eta + \text{const}$ , from where one can deduce that along the light ray there remains a constant product  $\omega a = \text{const}$ . A light ray with frequency  $\omega_0$  emitted at a distance  $\chi$  and observed at the origin ( $\chi = 0$ ) at time  $\eta$  should then have the frequency  $\omega = \omega_0 \frac{a(\eta - \chi)}{a(\eta)} \approx \omega_0(1 - \chi \frac{a'}{a})$ , that is, redshifted, if the universe expands ( $a' > 0$ ). The proper distance  $l$  to the source of light is  $l = \chi a$ . Thus the frequency shift  $z$  can be written as  $z \equiv \frac{\omega_0 - \omega}{\omega_0} = \frac{a'}{a^2} l \equiv Hl$ , where  $H$  is the so called Hubble constant,  $H = \frac{a'}{a^2} = \frac{1}{a} \frac{da}{dt}$ . The current value of the Hubble constant is  $H \approx \frac{1}{(13 \text{ bil. years})}$ .

Inserting  $\frac{a'}{a^2} = H$  into Friedmann's equations leads to  $\frac{1}{a^2} = H^2 - \frac{\kappa\mu}{3}$  for a closed universe, and to  $\frac{1}{a^2} = \frac{\kappa\mu}{3} - H^2$  for an open universe. For the critical density  $\mu_c$ , such that  $\frac{\kappa\mu_c}{3} = H^2$ , the universe is flat.

The current measurements show that the relative density  $\Omega = \frac{\mu}{\mu_c}$  is close to one with an error about few per cent (flatness problem). About 30% of it is "dark matter" and about 70% is "dark energy". The visible matter constitutes only about 3% of the density.

### Exercises

1. Consider a flat (Euclidean) isotropic universe with the metric  $ds^2 = dt^2 - b^2(t)(dx^2 + dy^2 + dz^2)$ :

- (a) Calculate the Christoffel symbols.  $[\Gamma_{tx}^x = \Gamma_{ty}^y = \Gamma_{tz}^z = \frac{b'}{b}, \Gamma_{xx}^t = \Gamma_{yy}^t = \Gamma_{zz}^t = bb']$
- (b) Calculate the Ricci tensor and the Ricci scalar.  $[R_t^t = -3\frac{b''}{b}, R_x^x = R_y^y = R_z^z = -\frac{b''}{b} - 2\frac{b'^2}{b^2}]$
- (c) Write down the  $t$  component of the Einstein's equations (with perfect fluid).  $[3\frac{b'^2}{b^2} = \kappa\epsilon]$
- (d) Write down the energy conservation equation  $\frac{dV}{V} = -\frac{d\epsilon}{(\epsilon+p)}$ .  $[3\ln(b) = -\int \frac{d\epsilon}{(\epsilon+p)}]$
- (e) Integrate the equations for a matter dominated universe ( $p = 0, \epsilon = \mu$ ).  $[\mu b^3 = \text{const}, b = \text{const} \cdot t^{2/3}]$
- (f) Integrate the equations for a radiation dominated universe ( $p = \frac{\epsilon}{3}$ ).  $[\epsilon b^4 = \text{const}, b = \text{const} \cdot t^{1/2}]$

2. Interpret the cosmological red shift  $\frac{\omega_0 - \omega}{\omega_0} = Hl$  ( $l$  is the distance to the "red shifted" galaxy) as a Doppler effect and calculate the velocity with which a galaxy appears to be moving relative to the observer.