note6 [October 2, 2023]

# The source of gravitation

In this section we shall establish that the source of gravitation is the stress-energymomentum tensor of matter.

In Newtonian gravitation—which should be the *slow motion weak field limit* of general relativity—the source of the gravitational field is the mass density of matter. However the mass density of matter cannot be the source of gravitation in general relativity as it is not a generally covariant quantity. Actually, the mass density is the zero-zero component of the stress-energy tensor of matter (which, being a tensor, is a generally covariant quantity). Therefore one might guess (correctly) that in general relativity the source of gravitation is the stress-energy tensor of matter.<sup>1</sup>

### Action in general relativity

The equation for the gravitational field shall be derived from the total action S of gravitation and matter,

$$S = S_M + S_G , (1)$$

where  $S_G$  is the action of the gravitation, to be determined later, and  $S_M$  is the action of matter,<sup>2</sup> to be taken—in a generally covariant form, of course—directly from special relativity.

Unlike electromagnetism, where there is an explicit coupling term between the electromagnetic field and charges,  $-\int d\Omega A_a j^a$ , in general relativity there are no coupling terms between gravitation and matter. Indeed, in general relativity gravitation couples to matter automatically through the metric tensor in the action of the matter.

Since the gravitational field is contained in the metric tensor,  $g_{ab}$ , of the time-space, the variational principle to determine the gravitational field created by a given distribution of matter should be formulated as

The actual gravitational field,  $g_{ab}$ , created by a given distribution of matter is that for which a small variation of the gravitational field,  $g_{ab} \rightarrow g_{ab} + \delta g_{ab}$ , leads to vanishing variation of the action,

$$\delta(S_g + S_M)\Big|_{g_{ab} \to g_{ab} + \delta g_{ab}} = 0.$$
 (2)

The source of the gravitational field is then determined by the variation of the matter action,  $S_M$ , over the metric tensor,

$$\delta S_M \Big|_{g_{ab} \to g_{ab} + \delta g_{ab}} . \tag{3}$$

We shall show that this variation is actually given by the stress-energy-momentum tensor of matter.

## Action of matter in general relativity

The action of matter in general relativity has the same form as in special relativity, only rewritten, if needed, in a generally covariant way. Particularly, one has to substitute

$$d\Omega \rightarrow \sqrt{-g}d\Omega$$
, (4)

$$\varphi^{,a} \rightarrow g^{ab}\varphi_{,b}$$
, (5)

$$A^b_{,a} \rightarrow A^b_{;a}. \tag{6}$$

For example, the action of the electromagnetic field becomes

$$-\frac{1}{16\pi} \int F^{ab} F_{ab} \sqrt{-g} d\Omega , \qquad (7)$$

<sup>&</sup>lt;sup>1</sup>It could also be the trace of the stress-energy tensor, that is also a covariant quantity. However that gives the second Nordström's theory of gravitation that was regected by experimental data.

<sup>&</sup>lt;sup>2</sup> "Matter" in general relativity is everything but the gravitational field (the latter being curvature of space-time).

note6 [October 2, 2023]

the coupling between the electromagnetic field and currents,

$$-\int A_a j^a \sqrt{-g} d\Omega , \qquad (8)$$

the action of a hypothetical massless scalar field  $\phi$ ,

$$-\int g^{ab}\varphi_a\varphi_b\sqrt{-g}d\Omega. (9)$$

### Metric variation of matter action and energy-momentum tensor

We shal show here that the variation of the matter action under variation of the metric tensor produces the stress-energy-momentum tensor of matter. The latter becomes the source of the gravitational field.

The action of matter is typically written in terms of the Lagrangian of the matter,  $\mathcal{L}$ ,

$$S_M = \int \mathcal{L}\sqrt{-g}d\Omega \,. \tag{10}$$

The variation of this action under a variation  $g^{ab} \rightarrow g^{ab} + \delta g^{ab}$  of the metric tensor can be written in terms of a symmetric tensor  $T_{ab}$ ,

$$\delta S_M \doteq \frac{1}{2} \int T_{ab} \delta g^{ab} \sqrt{-g} d\Omega = -\frac{1}{2} \int T^{ab} \delta g_{ab} \sqrt{-g} d\Omega,$$

 ${\rm where}^3$ 

$$\frac{1}{2}\sqrt{-g}T_{ab}\delta g^{ab} \doteq \delta(\sqrt{-g}\mathcal{L}). \tag{11}$$

The tensor  $T_{ab}$  must be the stress-energy-momentum tensor, since in a flat space it satisfies the conservation law  $T_{,b}^{ab}=0$ . Indeed, consider an infinitesimal coordinate transformation,

$$x^a \to x^{'a} = x^a + \epsilon^a \,. \tag{12}$$

The variation of the metric tensor under this transformation can be written as<sup>4</sup>

$$\delta g^{ab} = \epsilon^{a;b} + \epsilon^{b;a} , \delta g_{ab} = -\epsilon_{a;b} - \epsilon_{b;a} .$$
 (17)

The variation of the action then takes the form

$$\delta S = \int T_{ab} \epsilon^{a;b} \sqrt{-g} d\Omega \,. \tag{18}$$

Integrating by parts<sup>5</sup> gives

$$\delta S = -\int T_{a;b}^b \epsilon^a \sqrt{-g} d\Omega \,. \tag{19}$$

$$x_b^{\prime a} = \delta_b^a + \epsilon_b^a \,. \tag{13}$$

Substituting this into the transformation rule for the metric tensor,

$$g_{ab}(x) = x'^{c}_{,a} x'^{d}_{,b} g'_{cd}(x'), \qquad (14)$$

gives

$$g_{ab}(x) = \left(\delta_a^c + \epsilon_{,a}^c\right) \left(\delta_b^d + \epsilon_{,b}^d\right) \left(g_{cd}'(x) + g_{cd,e}\epsilon^e + O((\epsilon)^2)\right)$$
(15)

$$= g'_{ab} + g_{ad}\epsilon^d_{,b} + g_{bd}\epsilon^d_{,a} + g_{ab,e}\epsilon^e + O\left((\epsilon)^2\right) = g'_{ab} + \epsilon_{a;b} + \epsilon_{b;a} + O\left((\epsilon)^2\right). \tag{16}$$

From  $g_{ab}g^{bc} = \delta^c_a$  follows  $g_{ab}\delta g^{bc} = -\delta g_{ab}g^{bc}$  and therefore  $T_{ab}\delta g^{ab} = -T^{ab}\delta g_{ab}$ .

Differentiating  $x'^a = x^a + \epsilon^a$  gives

<sup>&</sup>lt;sup>5</sup>using  $(AB)_{;a} = A_{;a}B + AB_{;a}$  and  $\sqrt{-g}A_{;c}^c = (\sqrt{-g}A^c)_{,c}$  and that a total differential does not contribute to

note6 [October 2, 2023]

The variation of the action under coordinate transformation is zero. Thus the tensor  $T_b^a$  satisfies the equation

$$T_{\cdot h}^{ab} = 0, (20)$$

which in a flat space turns into the stress-energy-momentum conservation equation  $T^{ab}_{,b} = 0$ . One can thus assume that the tensor is proportional to the canonical stress-energy-momentum tensor. Direct calculations show that the proportionality factor is equal unity.

#### Exercises

1. The Lagrangian density of the electromagnetic field is given as

$$\mathcal{L} = -\frac{1}{16\pi} F_{ab} F^{ab}.$$

Calculate the corresponding stress-energy-momentum tensor,  $T_{ab}$ , using the "metric" definition.

$$\delta \left( \sqrt{-g} \mathcal{L} \right)_{g^{ab} \to g^{ab} + \delta g^{ab}} = \frac{1}{2} \sqrt{-g} T_{ab} \delta g^{ab}.$$

Answer:

$$T_{ab} = \frac{1}{16\pi} \left( F_{cd} F^{cd} g_{ab} - 4 F_{ac} F_b^{\ c} \right) .$$

2. Consider a scalar field  $\varphi$  whose action in special relativity is given as

$$S = \int d\Omega \left( \frac{1}{2} \varphi^{,a} \varphi_{,a} - \frac{1}{2} m^2 \varphi^2 \right) .$$

(a) Calculate its "translation-invariance" stress-energy-momentum tensor,

$$T_b^a = \frac{\partial \mathcal{L}}{\partial \varphi_a} \varphi_{,b} - \mathcal{L} \delta_b^a .$$

(b) Rewrite the action in a generally covariant form and calculate its "metric" stress-energy-momentum tensor,

$$\delta \left( \sqrt{-g} \mathcal{L} \right)_{g^{ab} \to g^{ab} + \delta g^{ab}} = \frac{1}{2} \sqrt{-g} T_{ab} \delta g^{ab}$$

Discuss the results.

3. Show that for an infinitesimally small coordinate transformation  $x^a \to x^a + \epsilon^a$  the variation of the metric tensor is given as

$$\delta g^{ab} = \epsilon^{a;b} + \epsilon^{b;a}$$
,  $\delta g_{ab} = -\epsilon_{a;b} - \epsilon_{b;a}$ .

4. Derive—by metric variation—the stress-energy-momentum tensor of a particle with mass m and argue that it is the correct tensor.

Hints:

- (a) The action is  $S = -m \int ds$ .
- (b) Its variation over  $\delta g_{ab}$  must look like  $\delta S = -\frac{1}{2} \int ds T^{ab} \delta g_{ab} \,.$
- 5. Argue that the equation  $j_{,a}^a=0$  represents a conservation law. What about  $j_{,a}^a=0$ ?
- 6. Argue that the equation  $T^{ab}_{,a}=0$  represents a conservation law. What about  $T^{ab}_{;a}=0$  ?
- 7. Prove eq. (19), that is, that

$$\int T^{ab} \epsilon_{a;b} \sqrt{-g} d\Omega = - \int T^{ab}_{;b} \epsilon_a \sqrt{-g} d\Omega.$$

note6 [October 2, 2023]

8. Argue that under an arbitrary infinitesimal coordinate transformation  $x^a \to x^a + \epsilon^a$  the variation of the action of the matter is given as

$$0 = \delta S_m = -\frac{1}{2} \int T^{ab} \delta g_{ab} \sqrt{-g} d\Omega = \int T^{ab}_{;b} \epsilon_a \sqrt{-g} d\Omega \,.$$

Does this imply that i)  $T^{ab}=0$  ? ii)  $T^{ab}_{;b}=0$  ?